

# *Unconventional superconductivity in double quantum dots*

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# Introduction



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# Symmetry of pair amplitude

- Consider pair amplitude:

$$F_{\alpha\beta}(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -i\langle T\psi_\alpha(\mathbf{r}_1, t_1)\psi_\beta(\mathbf{r}_2, t_2) \rangle$$

- Pauli principle:

$$F_{\alpha\beta}(\mathbf{r}_1, t_1; \mathbf{r}_2, t_2) = -F_{\beta\alpha}(\mathbf{r}_2, t_2; \mathbf{r}_1, t_1)$$

- Different ways to realize antisymmetry



# Classification of superconductivity

Frequency	Spin	Orbital	
+ (even)	- (singlet)	+ (even)	even-singlet
+ (even)	+ (triplet)	- (odd)	even-triplet
- (odd)	+ (triplet)	+ (even)	odd-triplet
- (odd)	- (singlet)	- (odd)	odd-singlet

- Even-singlet:  $s$ -wave BCS superconductivity
- Even-triplet: Superfluid  $^3\text{He}$ ,  $\text{Sr}_2\text{RuO}_4$ , Majorana nanowires
- Odd-triplet: SFS heterostructures
- Odd-singlet: ???



# Unconventional SC in quantum dots

- Can **unconventional** SC be induced in quantum dots?
- Interplay between **superconductivity**, strong **Coulomb** interaction and **nonequilibrium**
- Quantum dots offer **easy tunability** of their properties
- Quantum dot as tunable source of even-singlet, even-triplet, odd-triplet and odd-singlet correlations?
- Signatures in **transport** properties?

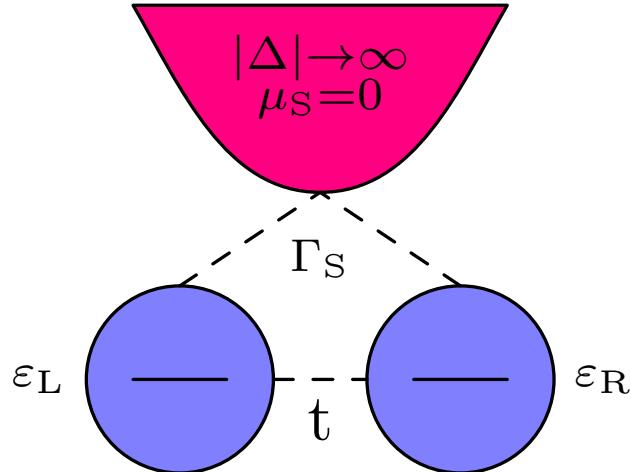


# Model



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# Model: Double dot



Effective Dot Hamiltonian for  $\Delta \rightarrow \infty$

Quantum dots  $i = L, R$

$$H_i = \varepsilon_i \sum_{\sigma} c_{i\sigma}^{\dagger} c_{i\sigma} + \mathbf{B}_i \cdot \hat{\mathbf{S}}_i$$

Intra- and interdot Coulomb interaction

$$H_{\text{inter}} = \sum_i U_i n_{i\uparrow} n_{i\downarrow} + U \sum_{\sigma\sigma'} n_{L\sigma} n_{R\sigma'}$$

Local and nonlocal proximity effect

$$H_{\text{prox}} = - \sum_i \frac{\Gamma_{Si}}{2} \left( c_{i\uparrow}^{\dagger} c_{i\downarrow}^{\dagger} + \text{H.c.} \right) - \frac{\Gamma_S}{2} \left( c_{R\uparrow}^{\dagger} c_{L\downarrow}^{\dagger} - c_{R\downarrow}^{\dagger} c_{L\uparrow}^{\dagger} + \text{H.c.} \right)$$

Interdot tunneling

$$H_{\text{tun}} = t \sum_{\sigma} (c_{L\sigma}^{\dagger} c_{R\sigma} + \text{H.c.})$$



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# Dynamics in Hilbert space

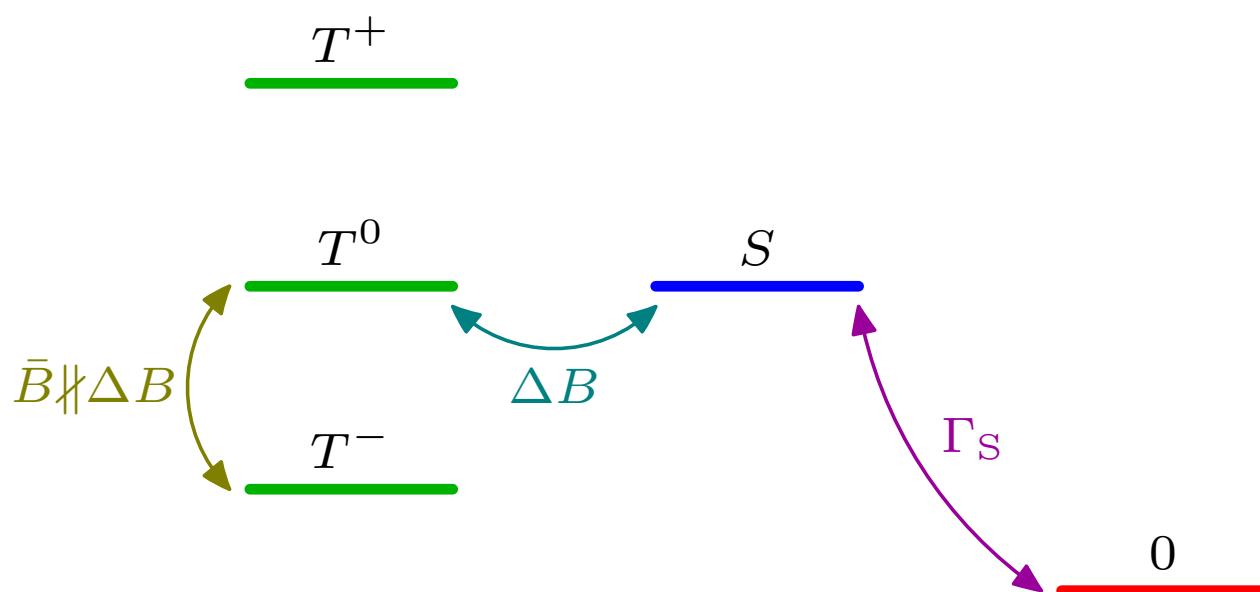
Relevant dot states with even occupation for  $U_i \rightarrow \infty$

Empty dot  $|0\rangle$ , nonlocal singlet  $|S\rangle$ , 3 triplets  $|T^\sigma\rangle$

Time evolution of density matrix elements  $P_{\chi_2}^{\chi_1} = \langle \chi_1 | \rho_{\text{dot}} | \chi_2 \rangle$

$$\frac{d}{dt} P_{\chi_2}^{\chi_1}(t) + i \sum_{\chi} (h_{\chi_1 \chi} P_{\chi_2}^{\chi} - h_{\chi \chi_2} P_{\chi}^{\chi_1}) (t) = 0$$

Matrix elements  $h_{\chi_1 \chi_2} = \langle \chi_1 | H_{\text{ddot}} | \chi_2 \rangle$



# Results



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# Pair amplitudes

Pair amplitude  $F_{i\sigma i'\sigma'}(t) = \langle T c_{i\sigma}(t) c_{i'\sigma'}(0) \rangle$

- Even/Odd-singlet pair amplitude:

$$F_{e/o}^S = (F_{L\downarrow R\uparrow} - F_{L\uparrow R\downarrow} \mp F_{R\uparrow L\downarrow} \pm F_{R\downarrow L\uparrow})/(2\sqrt{2})$$

- Even/Odd-triplet pair amplitudes:

$$F_{e/o}^{T^+} = (F_{L\uparrow R\uparrow} \mp F_{R\uparrow L\uparrow})/2$$

$$F_{e/o}^{T^0} = (F_{L\downarrow R\uparrow} + F_{L\uparrow R\downarrow} \mp F_{R\uparrow L\downarrow} \mp F_{R\downarrow L\uparrow})/(2\sqrt{2})$$

$$F_{e/o}^{T^-} = (F_{L\downarrow R\downarrow} \mp F_{R\downarrow L\downarrow})/2$$



# Even-frequency order parameters

Order parameter: Pair amplitude at equal times  $\Delta_e = F_e(0)$

For  $U_i \rightarrow \infty$ :

$$\Delta_e^S = P_0^S$$

$$\Delta_e^{T_\alpha} = P_0^{T_\alpha}$$

- Coupling to **superconductor**  $\Gamma_S$ : Even-singlet
- **Inhomogenous** magnetic field  $\Delta B_\alpha$ : Even-triplet along the inhomogeneity
- **Noncollinear** magnetic field  $\Delta \mathbf{B} \nparallel \bar{\mathbf{B}}$ : Even-triplet perpendicular to inhomogeneity

In agreement with findings for SFS heterostructures Bergeret, Volkov, Efetov, RMP 2005



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# Odd-frequency order parameters

- Pair amplitude at equal times **vanishes**  $F_o(0) = 0$
- Consider **time-derivative** of pair amplitude at equal times
$$\Delta_o = \left. \frac{dF_o(t)}{dt} \right|_{t=0}$$
- Express **odd-frequency** order parameter in terms of **even-frequency** order parameters and **local** and **nonlocal** expectation value of **charge** and **spin**

$$\Delta_o^T = -\frac{i}{2}\Delta\varepsilon\Delta_e^T + \frac{i}{2}\bar{\mathbf{B}}\Delta_e^S + \frac{1}{4}\Delta\mathbf{B} \times \Delta_e^T + \frac{i}{2\sqrt{2}} \left( \Gamma_S \mathbf{S} - \Gamma_{SL} \mathbf{S}_R^L - \Gamma_{SR} \mathbf{S}_L^R \right)$$

$$\Delta_o^S = -\frac{i}{2}\Delta\varepsilon\Delta_e^S + \frac{i}{2}\bar{\mathbf{B}} \cdot \Delta_e^T - \frac{i}{4\sqrt{2}} \left( \Gamma_S \Delta N + \Gamma_{SL} N_R^L - \Gamma_{SR} N_L^R \right)$$



# Transport signatures

How to probe unconventional correlations in transport?

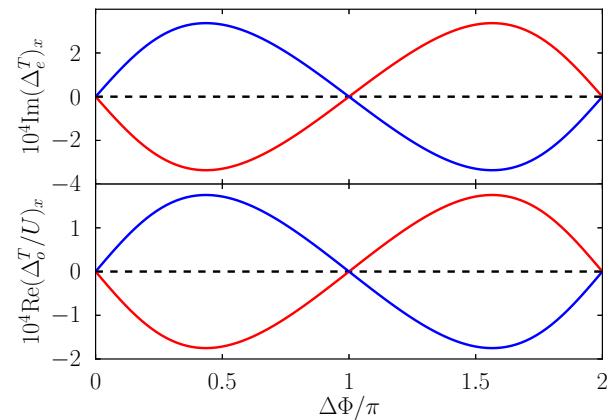
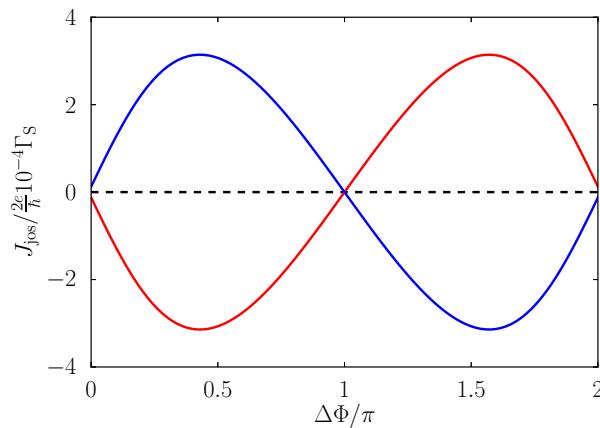
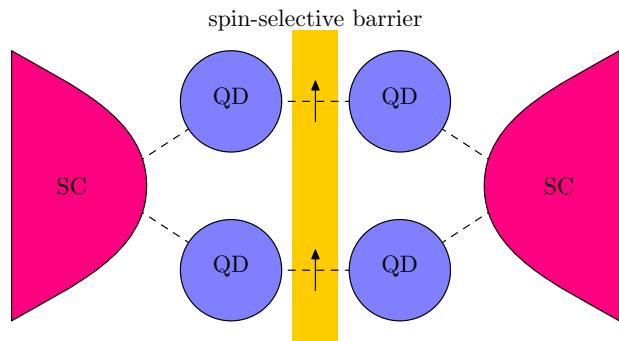
- Conventional SC only probes even-singlet pairing
- No unconventional electrodes available
- Indirect signatures needed!
- Contribution from even-singlet dominates transport
- Find situation where current due to even-singlet vanishes

Here: Two examples for detecting triplet correlations in Josephson current



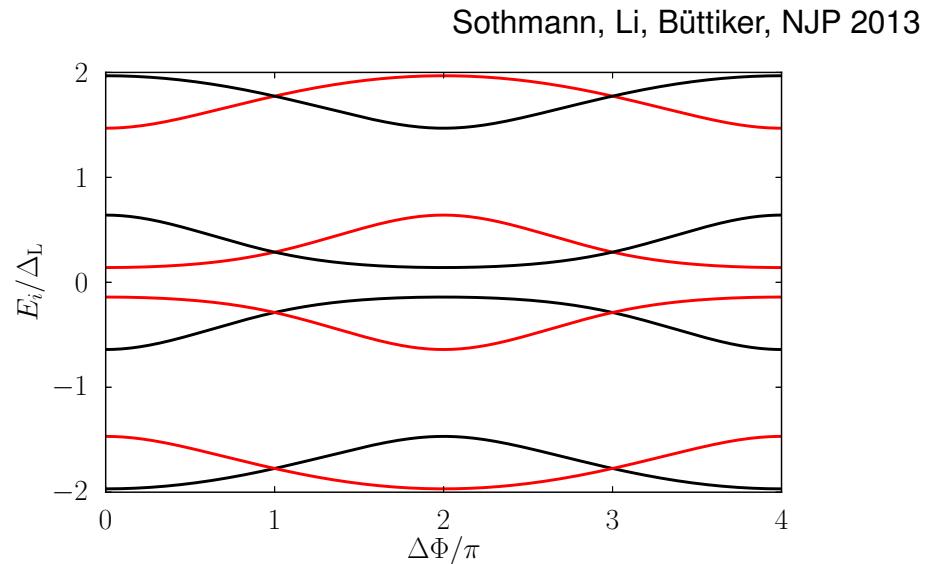
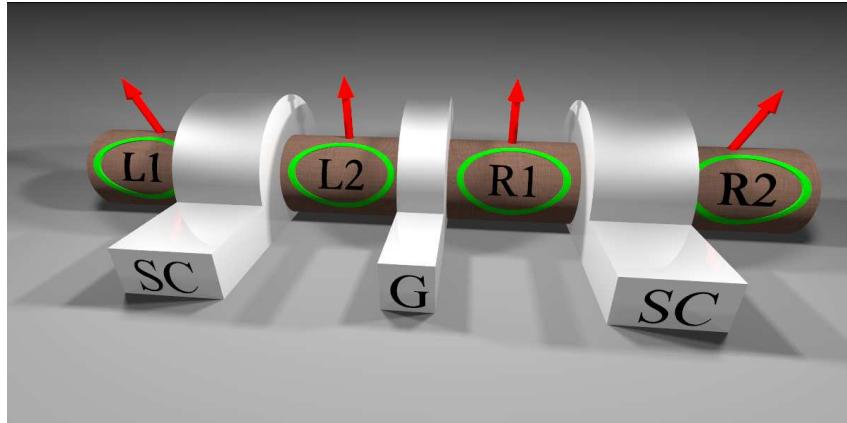
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# Triplet Josephson current



- Double dots connected via spin-selective barrier
- No Josephson current via spin singlet
- No magnetic field: Only even-singlet pairing: No Josephson current
- Inhomogenous magnetic field: Even- and odd-triplet pairing: Finite Josephson current
- $\pi$ -junction possible

# Fractional Josephson effect



- Quadruple quantum dot with inhomogenous magnetic field
- Dots can be tuned to host Majorana fermions
- Fractional Josephson effect possible



# Summary



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# Summary & Outlook

- Quantum dots can host **unconventional** SC
- **All types** of unconventional SC can be generated
- **Tune** system between different types of unconventional SC
- Signatures of unconventional SC in **transport**
- Can one find clear signatures of odd-singlet SC in this system?

